

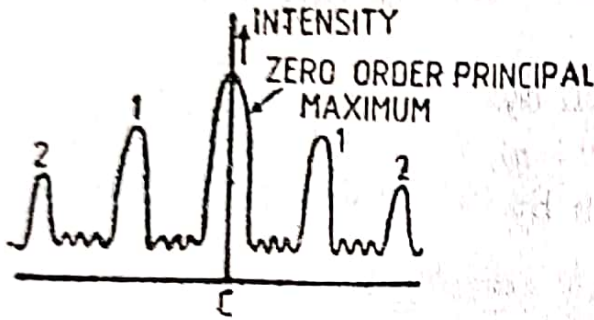
7.7. Fraunhofer diffraction at N slits

$S_1, S_2, \dots, S_n \rightarrow N$ parallel slits, each of width ' a ' and separated by opaque space ' b '.

$L \rightarrow$ a collimating lens.

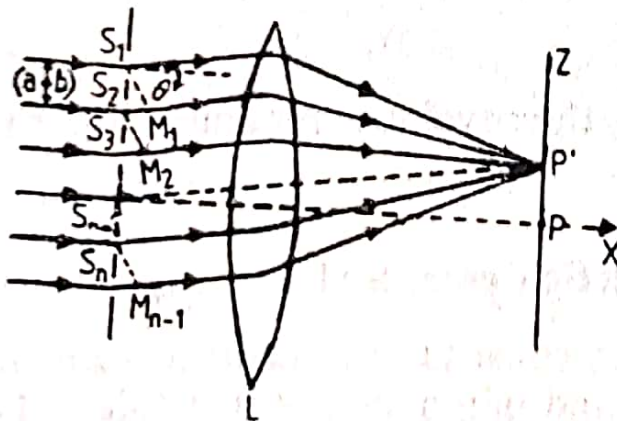
$Z \rightarrow$ screen.

Result : The pattern obtained on the screen is called the Fraunhofer diffraction pattern due to N parallel equidistant slits. The pattern consists of a central maximum called Zero order principal maximum followed by a number of principal maxima on either side. There are $(N-1)$ equi-spaced minima between any two consecutive principal maxima.



There are $(N-2)$ other maxima coming alternatively with the minima between two consecutive principal maxima. These maxima are called secondary maxima. The intensity distribution is shown in the figure.

Theory : A plane wavefront of monochromatic light is incident normally on N parallel slits.



The light is diffracted through N slits and is focussed by a lens L on the screen Z placed in the focal plane of L . The wave diffracted from all the slits in direction θ are equivalent to N parallel waves, each wave starting from mid-points, S_1, S_2, \dots, S_n of the slits.

Let $S_1M_1, S_2M_2, \dots, S_{n-1}M_{n-1}$ be perpendiculars dropped from S_1, S_2, \dots, S_{n-1} on to the diffracted rays. The path difference between the waves from

S_1 and S_2 is

S_2 and S_3 is

S_{n-1} and S_n is

$$S_2M_1 = (a+b)\sin\theta = d\sin\theta,$$

$$S_3M_2 = d\sin\theta,$$

$$S_nM_{n-1} = d\sin\theta,$$

Thus as we pass from one wave to another, the path goes on increasing by the same amount $d\sin\theta$.

To obtain an expression for the intensity distribution due to diffraction at N slits, the expression for instantaneous displacement (dZ) has to be integrated for N slits.

For a single slit,

$$Z = k \int_{-a/2}^{a/2} \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{y \sin \theta}{\lambda} \right) dy = k \int_{-a/2}^{a/2} \varphi(y) dy$$

when $\varphi(y) = \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{y \sin \theta}{\lambda} \right)$.

For N slits :

$$\begin{aligned} Z &= k \int_{-a/2}^{a/2} \varphi(y) dy + \int_{d-a/2}^{d+a/2} \varphi(y) dy \\ &\quad + \int_{2d-a/2}^{2d+a/2} \varphi(y) dy + \dots + \int_{(N-1)d-a/2}^{(N-1)d+a/2} \varphi(y) dy \\ &= ka \frac{\sin \alpha}{\alpha} \left[\sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} \right) + \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{d \sin \theta}{\lambda} \right) \right. \\ &\quad \left. + \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{2d \sin \theta}{\lambda} \right) + \dots \right. \\ &\quad \left. + \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{(N-1) d \sin \theta}{\lambda} \right) \right] \\ &= ka \frac{\sin \alpha}{\alpha} \cdot \frac{\sin N\beta}{\sin \beta} \cdot \sin 2\pi \left(\frac{t}{T} - \frac{r}{\lambda} + \frac{(N-1) d \sin \theta}{2\lambda} \right) \end{aligned}$$

where $\alpha = \pi a \sin \theta / \lambda$, and $\beta = \pi d \sin \theta / \lambda$.

The intensity at a point P' is given by

$$I = k^2 a^2 \left(\frac{\sin^2 \alpha}{\alpha^2} \right) \left(\frac{\sin^2 N\beta}{\sin^2 \beta} \right) = I_0 \frac{\sin^2 \alpha}{\alpha^2} \cdot \frac{\sin^2 N\beta}{\sin^2 \beta} \left[\because I_0 = k^2 a^2 \right]$$

The factor $(\sin^2 \alpha / \alpha^2)$ represents the diffraction pattern due to a single slit. The factor $(\sin^2 N\beta / \sin^2 \beta)$ represents the interference effects due to the secondary waves from the N slits.

7.8. (i) Principal Maxima

$$\lim_{\beta \rightarrow 0} \frac{\sin N\beta}{\sin \beta} = N.$$

When $\beta \rightarrow 0$, the resultant intensity of maxima is

$$I = I_0 \frac{\sin^2 \alpha}{\alpha^2} \cdot N^2 = I_0 N^2 \frac{\sin^2 \alpha}{\alpha^2}.$$

Thus if we increase the number of slits, the intensity of principal maxima increases. The directions of principal maxima are given by

$$\begin{aligned} \sin \beta &= 0 \\ \Rightarrow \beta &= \pm n\pi \text{ where } n = 0, 1, 2, \dots \end{aligned}$$

$$\Rightarrow \frac{\pi d \sin \theta}{\lambda} = \pm n\pi$$

$$\Rightarrow d \sin \theta = (a+b) \sin \theta = \pm n\lambda.$$

When $n=0$, we get $\theta=0$. This is the direction in which all the waves arrive in the same phase and hence we get a bright central maximum. This is called the Zero order principal maximum. By putting $n=1, 2, 3, \dots$ we obtain first, second, third... order principal maxima respectively.

(ii) Minima : For minima, we must have

$$\sin N\beta = 0 \quad (\because \sin \beta \neq 0)$$

$$\Rightarrow N\beta = \pm m\pi \Rightarrow N \frac{\pi d \sin \theta}{\lambda} = \pm m\pi$$

$$\Rightarrow N(a+b) \sin \theta = \pm m\lambda.$$

Here m can have all integral values except $0, N, 2N, 3N, \dots$ which give $\sin \beta = 0$ which is the condition for principal maxima. The positive and negative values of m indicate that the minima of a given order lie symmetrically on either side of the central principal maximum.

N.B. $m=0$ gives principal maximum of zero order. $m=1, 2, 3, \dots (N-1)$ give $(N-1)$ minima, $m=N$ gives principal maximum of first order.

(iii) Secondary Maxima

Between two consecutive principal maxima there are $(N-1)$ minima and hence $(N-2)$ maxima. These maxima are called secondary maxima. Their positions are obtained by

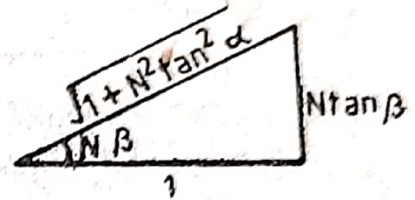
$$\Rightarrow I_0 \frac{\sin^2 \alpha}{\alpha^2} \cdot 2 \left(\frac{\sin N\beta}{\sin \beta} \right) \cdot \frac{\sin \beta \cdot N \cos N\beta - \sin N\beta \cdot \cos \beta}{\sin^2 \beta} = 0$$

$$\Rightarrow N \cos N\beta \cdot \sin \beta - \sin N\beta \cdot \cos \beta = 0$$

$$\Rightarrow \tan N\beta = N \tan \beta.$$

The roots of this equation (other than $\beta = \pm n\pi$) give the position of secondary maxima.

In order to find the intensity of secondary maxima, we have



$$\sin N\beta = N \tan \beta / \sqrt{1 + N^2 \tan^2 \beta}$$

$$\Rightarrow \sin^2 N\beta = N^2 \tan^2 \beta / (1 + N^2 \tan^2 \beta)$$

$$\Rightarrow \frac{\sin^2 N\beta}{\sin^2 \beta} = \frac{N^2 \tan^2 \beta}{(1 + N^2 \tan^2 \beta) \sin^2 \beta} = \frac{N^2}{1 + (N^2 - 1) \sin^2 \beta}$$

Hence the intensity of secondary maxima is given by

$$I_s = I_0 \frac{\sin^2 \alpha}{\alpha^2} \cdot \frac{N^2}{1 + (N^2 - 1) \sin^2 \beta}$$

Thus as N increases the intensity of secondary maxima decreases. When N is very large (Plane grating), the secondary maxima are not visible. There is uniform darkness between any two consecutive principal maxima.